

## Superconductivity in SrNi<sub>2</sub>As<sub>2</sub> single crystals

E. D. Bauer, F. Ronning, B. L. Scott, and J. D. Thompson  
 Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA  
 (Received 6 August 2008; published 7 November 2008)

The electrical resistivity  $\rho(T)$  and heat capacity  $C(T)$  on single crystals of SrNi<sub>2</sub>As<sub>2</sub> and EuNi<sub>2</sub>As<sub>2</sub> are reported. While there is no evidence for a structural transition in either compound, SrNi<sub>2</sub>As<sub>2</sub> is found to be a bulk superconductor at  $T_c=0.62$  K with a Sommerfeld coefficient of  $\gamma=8.7$  mJ/mol K<sup>2</sup> and a small upper critical field  $H_{c2}\sim 200$  Oe. No superconductivity was found in EuNi<sub>2</sub>As<sub>2</sub> above 0.4 K, but anomalies in  $\rho$  and  $C$  reveal that magnetic order associated with the Eu<sup>2+</sup> magnetic moments occurs at  $T_m=14$  K.

DOI: 10.1103/PhysRevB.78.172504

PACS number(s): 74.10.+v, 74.25.Bt, 74.70.Dd

The ThCr<sub>2</sub>Si<sub>2</sub> structure type is well known for accommodating a superconducting ground state, particularly in the heavy fermion community with superconductors such as CeCu<sub>2</sub>Si<sub>2</sub> and URu<sub>2</sub>Si<sub>2</sub>.<sup>1</sup> Soon after the discovery of superconductivity in LaFeAsO<sub>1-x</sub>F<sub>x</sub> at  $T_c=26$  K with the structure type ZrCuSiAs,<sup>2</sup> it was realized that the related compounds (AFe<sub>2</sub>As<sub>2</sub>, with A=Ba, Sr, Ca, and Eu) in the ThCr<sub>2</sub>Si<sub>2</sub> structure are also superconducting either with doping<sup>3-7</sup> or under pressure.<sup>8-10</sup> While systems with Fe<sub>2</sub>As<sub>2</sub> planes have the highest  $T_c$ 's to date, superconductivity has been found in both structure types with either Ni<sub>2</sub>P<sub>2</sub> (Refs. 11 and 12) or Ni<sub>2</sub>As<sub>2</sub> (Refs. 13-16) layers.

Here we report the observation of superconductivity in single crystals of SrNi<sub>2</sub>As<sub>2</sub> at  $T_c=0.62$  K, as determined by heat capacity, in the absence of a structural phase transition (below 400 K). Following our initial observation of superconductivity in BaNi<sub>2</sub>As<sub>2</sub>,<sup>16</sup> this represents the second superconducting system in the ThCr<sub>2</sub>Si<sub>2</sub> structure with Ni<sub>2</sub>As<sub>2</sub> layers. In addition, we report that our EuNi<sub>2</sub>As<sub>2</sub> single crystals grown from Pb flux are not superconducting above 0.4 K.

Single crystals of SrNi<sub>2</sub>As<sub>2</sub> and EuNi<sub>2</sub>As<sub>2</sub> were grown in Pb flux in the ratio (Sr, Eu):Ni:As:Pb=1:2:2:20. The starting elements were placed in an alumina crucible and sealed under vacuum in a quartz ampoule. The ampoule was placed in a furnace and slowly heated to 1050 °C, as described in Ref. 16. The sample was then cooled slowly (5 °C hr<sup>-1</sup>) to 600 °C, at which point the excess Pb flux was removed with the aid of a centrifuge. For SrNi<sub>2</sub>As<sub>2</sub> the resulting platelike crystals were heavily embedded in a yet unidentified needle-like impurity phase. From single-crystal x-ray refinements, the platelike samples were confirmed to crystallize in the ThCr<sub>2</sub>Si<sub>2</sub> tetragonal structure (space-group 139, *I4/mmm*). The refinement for SrNi<sub>2</sub>As<sub>2</sub> [ $R(I>2\sigma)=3.7\%$ ] at 124 K yields lattice parameters  $a=4.1374(8)$  Å and  $c=10.188(4)$  Å and fully occupied (>98%) atomic positions Sr 2a(0,0,0), Ni 4d(0.5,0,0.25), and As 4e(0,0,z) with  $z=0.3634(1)$  consistent with previous reports.<sup>17-19</sup> The refinement for EuNi<sub>2</sub>As<sub>2</sub> [ $R(I>2\sigma)=5.09\%$ ] at 124 K gives lattice parameters  $a=4.0964(6)$  Å and  $c=10.029(3)$  Å and fully occupied (>98%) atomic positions Eu 2a(0,0,0), Ni 4d(0.5,0,0.25), and As 4e(0,0,z) with  $z=0.3674(2)$  also consistent with previous reports.<sup>20,21</sup>

The in-plane electrical resistivity data for ANi<sub>2</sub>As<sub>2</sub> (A=Ba, Sr, Eu) is shown in Fig. 1. All samples exhibit metallic behavior. SrNi<sub>2</sub>As<sub>2</sub> is a relatively good metal with a RRR [ $=\rho(300\text{ K})/\rho(4\text{ K})$ ] of 11, and a residual resistivity of

7  $\mu\Omega$  cm. The resistivity of EuNi<sub>2</sub>As<sub>2</sub> exhibits a kink at  $T_m=14$  K associated with magnetic ordering of the Eu<sup>2+</sup> moments, consistent with previous reports.<sup>21</sup> For SrNi<sub>2</sub>As<sub>2</sub>, there is no evidence of a structural transition below 400 K, in contrast to BaNi<sub>2</sub>As<sub>2</sub>, which has a clear first-order transition at  $T_0=130$  K. The lack of a phase transition in SrNi<sub>2</sub>As<sub>2</sub> is also provided by the heat-capacity data shown in the inset of Fig. 3.<sup>22</sup>

Figure 2 presents the low-temperature in-plane resistivity data for SrNi<sub>2</sub>As<sub>2</sub> with fields applied parallel and perpendicular to the *c* axis. In zero field, a sharp superconducting transition is observed at  $T_c=0.66$  K, defined as the midpoint of the resistive anomaly. With increasing magnetic field, the transition remains sharp and is quickly suppressed. The specific heat shown in Fig. 3 confirms the bulk nature of superconductivity in SrNi<sub>2</sub>As<sub>2</sub>. The zero resistance state coincides exactly with the onset of the specific-heat transition, from which we extract a superconducting transition temperature of  $T_c=0.62$  K by an equal area construction. From a fit to the data from 0.7 to 3 K of  $C/T=\gamma+\beta T^2+\delta T^4$ , a Sommerfeld coefficient  $\gamma=8.7$  mJ/mol K<sup>2</sup> is obtained. Using this value, the ratio of the specific-heat jump at  $T_c$  to the electronic specific heat is estimated to be  $\Delta C/\gamma T_c\approx 1.0$ . From the

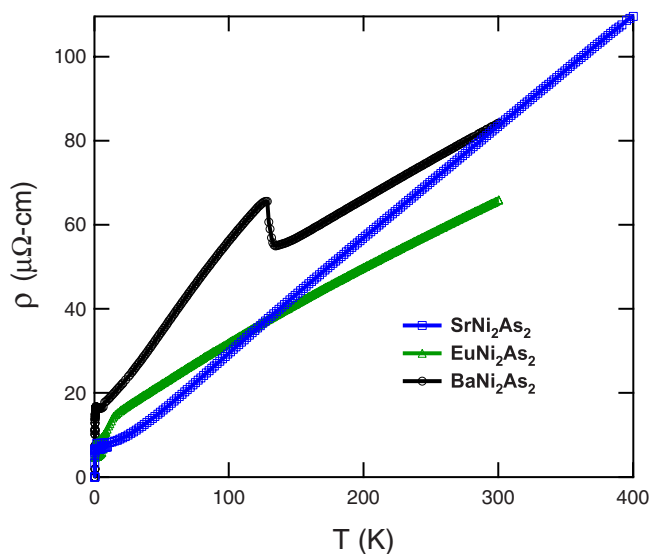


FIG. 1. (Color online) In-plane electrical resistivity  $\rho(T)$  ( $I\parallel ab$ ) for selected ANi<sub>2</sub>As<sub>2</sub> (A=Ba, Sr, Eu) compounds.

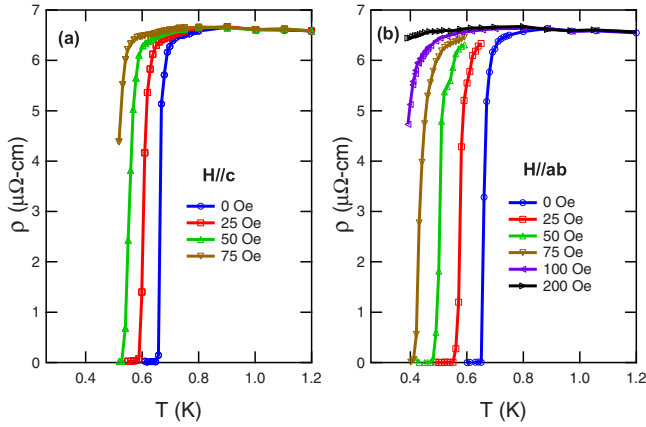


FIG. 2. (Color online) Electrical resistivity  $\rho(T)$  of  $\text{SrNi}_2\text{As}_2$  showing the superconducting transition for (a)  $H\parallel c$  and (b)  $H\parallel ab$ . The current was maintained perpendicular to the magnetic field.

$\beta$  coefficient =  $0.67 \text{ mJ/mol K}^4$ , one obtains a Debye temperature  $\Theta_D = 244 \text{ K}$ .

The magnetic field-temperature  $H$ - $T$  phase diagram of  $\text{SrNi}_2\text{As}_2$  is shown in Fig. 4 determined from the  $\rho(T)$  curves in Fig. 2, along with the data for  $\text{BaNi}_2\text{As}_2$  for comparison.<sup>16</sup> For  $\text{SrNi}_2\text{As}_2$ , the zero-temperature orbital critical field<sup>23</sup>  $H_{c2}^*(0) = -0.7T_c dH_{c2}/dT_c$  is determined to be 210 and 150 Oe for  $H\parallel c$  and  $H\parallel ab$ , respectively. From this, the superconducting coherence length is estimated via  $H_{c2}^*(0) = \Phi_0/2\pi\xi^2$ ,<sup>24</sup> yielding  $\xi^{ab} = 1477 \text{ \AA}$  and  $\xi^c = 1250 \text{ \AA}$ . Assuming  $\text{SrNi}_2\text{As}_2$  is in the clean limit, values for the Fermi velocity  $v_F^{ab} = 7.1 \times 10^6 \text{ cm/s}$  and  $v_F^c = 6.0 \times 10^6 \text{ cm/s}$  are obtained from  $\xi_0 = 0.18\hbar v_F/k_B T_c$ . Surprisingly, while the absolute value of  $T_c$  in zero magnetic field is very similar for the two compounds, the anisotropy ( $H_{c2}^{ab}/H_{c2}^c$ ) of the upper critical field, which is a factor of 2.1 in  $\text{BaNi}_2\text{As}_2$ , reverses

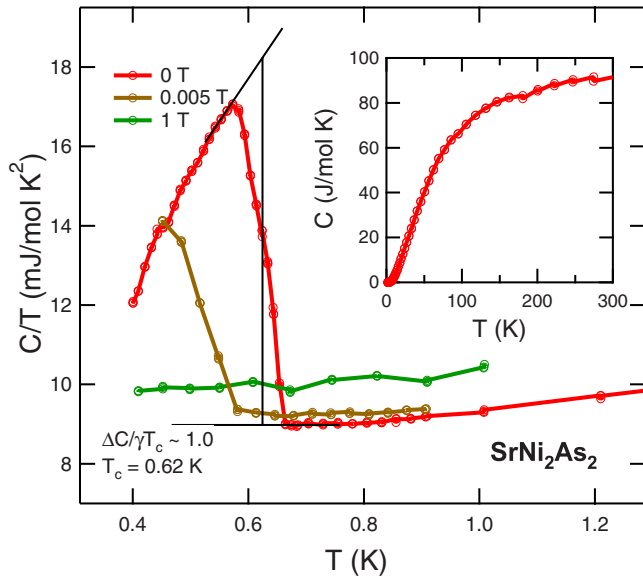


FIG. 3. (Color online) Low-temperature specific heat  $C$  versus temperature  $T$  of  $\text{SrNi}_2\text{As}_2$  (Ref. 22) for various magnetic fields ( $H\parallel c$ ). The inset displays the high-temperature heat capacity with no indication of a first-order phase transition at higher temperatures.

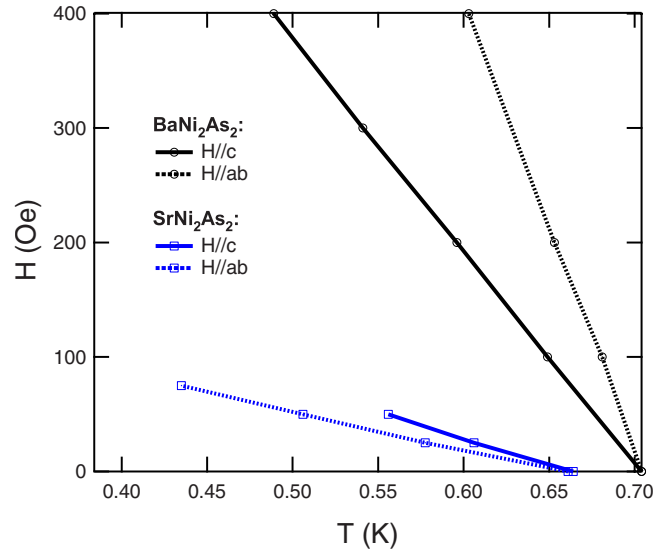


FIG. 4. (Color online) Magnetic field-temperature  $H$ - $T$  phase diagram of  $\text{SrNi}_2\text{As}_2$ . The upper critical field  $H_{c2}$  for  $H\parallel c$  and  $H\parallel ab$  was determined by the resistive midpoint in Figs. 2(a) and 2(b). The data for  $\text{BaNi}_2\text{As}_2$  from Ref. 16 is for the zero resistance state.

sign ( $H_{c2}^c/H_{c2}^{ab} = 1.4$ ) in  $\text{SrNi}_2\text{As}_2$  and the overall magnitude of the upper critical field is nearly an order of magnitude smaller. The differences in the  $H_{c2}$  between the two compounds may be due to changes in electronic structure resulting from the structural transition in  $\text{BaNi}_2\text{As}_2$  that is not present in  $\text{SrNi}_2\text{As}_2$ .

The specific heat, plotted as  $C/T$ , for  $\text{EuNi}_2\text{As}_2$  is shown in Fig. 5 in magnetic fields up to 9 T ( $H\parallel c$ ). A sharp anomaly occurs at the magnetic ordering temperature  $T_m = 14 \text{ K}$  [consistent with the kink in  $\rho(T)$ , Fig. 1], as well as a broader hump at  $\sim 4 \text{ K}$ . A magnetic field along the  $c$  axis modestly suppresses the transition, consistent with previous reports of antiferromagnetic ordering in polycrystalline samples.<sup>21,25</sup>

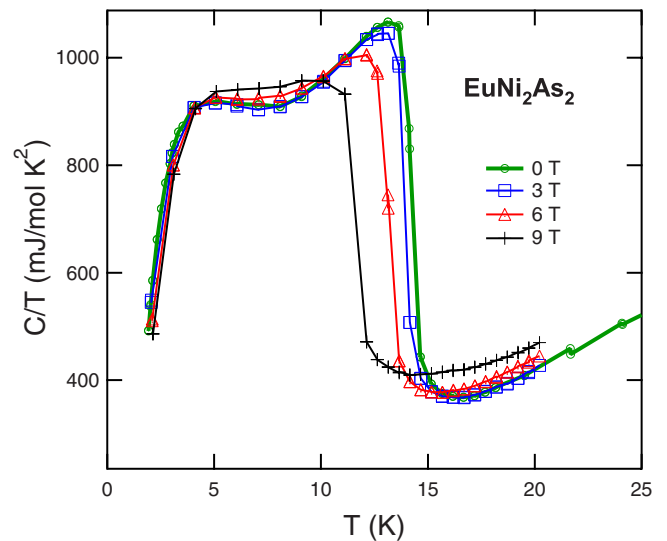


FIG. 5. (Color online) Heat-capacity data versus temperature for  $\text{EuNi}_2\text{As}_2$  in zero and applied magnetic field (Ref. 22). The magnetic field was applied along the  $c$  axis.

There is no indication of superconductivity above 0.4 K in  $\text{EuNi}_2\text{As}_2$ .

It is interesting that superconductivity with very similar transition temperatures is found both in  $\text{BaNi}_2\text{As}_2$  (Ref. 16) and in  $\text{SrNi}_2\text{As}_2$  despite the differences in structural parameters caused by the smaller  $\text{Sr}^{2+}$  ions, as well as the presence of a first-order structural transition in  $\text{BaNi}_2\text{As}_2$  that is possibly also magnetic. Recent theoretical work by Subedi *et al.*<sup>26</sup> indicate that the superconducting properties of the related Ni-analog  $\text{LaNiPO}$  may be explained within a conventional electron-phonon approach yielding a low value of  $T_c = 2.6\text{K}$ , consistent with experiment; the authors go on to suggest that the Fe-As superconductors may be in a separate class from their Ni-based counterparts. However, a scenario was put forth by Cvetkovic and Tesanovic;<sup>27</sup> involving a multiband Fermi surface in layered FeAs superconductors to

produce the large values of  $T_c$  may also be an appropriate description of these  $\text{ANi}_2\text{As}_2$  superconductors as well. Further work is in progress to elucidate the nature of the superconductivity in these Ni-based materials and its relation to fine details of the electronic structure.

In conclusion, specific-heat and electrical resistivity measurements of  $\text{SrNi}_2\text{As}_2$  single crystals reveal bulk superconductivity at 0.62 K, which shows no sign of a structural and/or magnetic anomaly below 400 K. Magnetic ordering associated with the Eu magnetic moments is observed in single crystalline  $\text{EuNi}_2\text{As}_2$ . No evidence for superconductivity is observed in this compound above 0.4 K.

Work at Los Alamos National Laboratory was performed under the auspices of the U.S. Department of Energy.

- 
- <sup>1</sup>D. W. Hess, P. S. Riseborough, and J. L. Smith, *Encyclopedia of Applied Physics* (VCH, New York, 1993), Vol. 7, p. 435.
- <sup>2</sup>Y. Kamihara, T. Watanabe, M. Hirano, and H. Hosono, *J. Am. Chem. Soc.* **130**, 3296 (2008).
- <sup>3</sup>M. Rotter, M. Tegel, and D. Johrendt, *Phys. Rev. Lett.* **101**, 107006 (2008).
- <sup>4</sup>G. F. Chen, Z. Li, G. Li, W. Z. Hu, J. Dong, X. D. Zhang, P. Zheng, N. L. Wang, and J. L. Luo, *Chin. Phys. Lett.* **25**, 3403 (2008).
- <sup>5</sup>K. Sasmal, B. Lv, B. Lorenz, A. Guloy, F. Chen, Y. Xue, and C. W. Chu, *Phys. Rev. Lett.* **101**, 107007 (2008).
- <sup>6</sup>G. Wu, H. Chen, T. Wu, Y. L. Xie, Y. J. Yan, R. H. Liu, X. F. Wang, J. J. Ying, and X. H. Chen, *J. Phys.: Condens. Matter* **20**, 422201 (2008).
- <sup>7</sup>H. S. Jeevan, Z. Hossain, C. Geibel, and P. Gegenwart, *Phys. Rev. B* **78**, 092406 (2008).
- <sup>8</sup>T. Park, E. Park, H. Lee, T. Klimczuk, E. D. Bauer, F. Ronning, and J. D. Thompson, *J. Phys.: Condens. Matter* **20**, 322204 (2008).
- <sup>9</sup>M. S. Torikachvili, S. L. Budko, N. Ni, and P. C. Canfield, *Phys. Rev. Lett.* **101**, 057006 (2008).
- <sup>10</sup>P. L. Alireza, J. Gillet, Y. T. Chris Ko, S. E. Sebastian, and G. G. Lonzarich, arXiv:0807.1896, *J. Phys.: Condens. Matter* (to be published).
- <sup>11</sup>T. Watanabe, H. Yanagi, T. Kamiya, Y. Kamihara, H. Hiramatsu, M. Hirano, and H. Hosono, *Inorg. Chem.* **46**, 7719 (2007).
- <sup>12</sup>T. Mine, H. Yanagi, T. Kamiya, Y. Kamihara, M. Hirano, and H. Hosono, *Solid State Commun.* **147**, 111 (2008).
- <sup>13</sup>T. Watanabe, H. Yanagi, Y. Kamihara, T. Kamiya, M. Hirano, and H. Hosono, *J. Solid State Chem.* **181**, 2117 (2008).
- <sup>14</sup>L. Fang, H. Yang, P. Cheng, X. Zhu, G. Mu, and H.-H. Wen, *Phys. Rev. B* **78**, 104528 (2008).
- <sup>15</sup>Z. Li, G. F. Chen, J. Dong, G. Li, W. Z. Hu, D. Wu, S. K. Su, P. Zheng, T. Xiang, N. L. Wang, and J. L. Luo, *Phys. Rev. B* **78**, 060504(R) (2008).
- <sup>16</sup>F. Ronning, N. Kurita, E. D. Bauer, B. L. Scott, T. Park, T. Klimczuk, R. Movshovich, and J. D. Thompson, *J. Phys. Condens. Matter* **20**, 342203 (2008).
- <sup>17</sup>A. Mewis and A. Distler, *Z. Naturforsch. B* **35**, 391 (1980).
- <sup>18</sup>M. Pfisterer and G. Nagorsen, *Z. Naturforsch. B* **35**, 703 (1980).
- <sup>19</sup>M. Pfisterer and G. Nagorsen, *Z. Naturforsch. B* **38**, 811 (1983).
- <sup>20</sup>W. Jeitschko, W. K. Hofmann, and L. J. Terbüchte, *J. Less-Common Met.* **137**, 133 (1988).
- <sup>21</sup>E. H. El Ghadraoui, J. Y. Pivan, R. Guérin, O. Pena, J. Padiou, and M. Sergent, *Mater. Res. Bull.* **23**, 1345 (1988).
- <sup>22</sup>Due to the small size of the heat-capacity signal, we are unable to properly subtract the addenda. This is particularly evident at high temperatures where we do not recover the Dulong-Petit value. We believe this is also responsible for the shift of approximately 1 mJ/mol K<sup>2</sup> at low temperatures and a magnetic field of 1 T.
- <sup>23</sup>N. R. Werthamer, E. Helfand, and P. C. Hohenberg, *Phys. Rev.* **147**, 295 (1966).
- <sup>24</sup>See, for example, M. Tinkham, *Introduction to Superconductivity* (McGraw-Hill, New York, 1975).
- <sup>25</sup>H. Raffius, E. Mörsen, B. D. Mosel, W. Müller-Warmuth, W. Jeitschko, L. Terbüchte, and T. Vomhof, *J. Phys. Chem. Solids* **54**, 135 (1993).
- <sup>26</sup>A. Subedi, D. J. Singh, and M.-H. Du, *Phys. Rev. B* **78**, 060506 (2008).
- <sup>27</sup>V. Cvetkovic and Z. Tesanovic, arXiv:0804.4678 (unpublished).